# Black Holes, Wormholes & Entanglement

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#### Agenda

Entanglement

Einstein's Field Equation

Black Holes

Properties of Black Holes

ER = EPR

Summary

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 $\mathbf{ER} = \mathbf{EPR}$ 

Summary

### Definition

 $|\varphi\rangle \in \mathbb{H}_1 \otimes \cdots \otimes \mathbb{H}_n$  is separable : $\Leftrightarrow$ 

 $|\varphi\rangle = |\psi_1\rangle \otimes \cdots \otimes |\psi_n\rangle$  with  $|\psi_i\rangle \in \mathbb{H}_i$  for  $1 \le i \le n$ 

 $|\varphi\rangle$  is *entangled* : $\Leftrightarrow$   $|\varphi\rangle$  is not separable

# Examples

$$\begin{split} \left| \varphi^{+} \right\rangle &= \frac{1}{\sqrt{2}} (\left| 0 \right\rangle \otimes \left| 0 \right\rangle + \left| 1 \right\rangle \otimes \left| 1 \right\rangle) \\ \left| \varphi^{-} \right\rangle &= \frac{1}{\sqrt{2}} (\left| 0 \right\rangle \otimes \left| 0 \right\rangle - \left| 1 \right\rangle \otimes \left| 1 \right\rangle) \\ \left| \psi^{+} \right\rangle &= \frac{1}{\sqrt{2}} (\left| 0 \right\rangle \otimes \left| 1 \right\rangle + \left| 1 \right\rangle \otimes \left| 0 \right\rangle) \\ \left| \psi^{-} \right\rangle &= \frac{1}{\sqrt{2}} (\left| 0 \right\rangle \otimes \left| 1 \right\rangle - \left| 1 \right\rangle \otimes \left| 0 \right\rangle) \end{split}$$

$$|W_n\rangle = \frac{1}{\sqrt{n}} (|0...001\rangle + |0...010\rangle + |0...100\rangle + ...+ |1...000\rangle)$$

$$|\Psi_{GHZ}^n\rangle = \frac{1}{\sqrt{2}} \left(|0\cdots0\rangle + |1\cdots1\rangle\right)$$

### A Phenomenon

$$\frac{1}{\sqrt{2}} (|00\rangle + |01\rangle) = |0\rangle \otimes \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$
(separabel)
$$\frac{1}{\sqrt{2}} (|00\rangle + |01\rangle)$$

Measuring the first qbit results in  $|0\rangle$  with probability 1. The second qbit will be measured as  $|0\rangle$  or  $|1\rangle$  with probability 1/2 This state is *entangled* 

$$\frac{1}{\sqrt{2}} \left( \left| 00 \right\rangle + \left| 11 \right\rangle \right)$$

Measuring the first qbit results in  $|0\rangle$  or  $|1\rangle$  with equal probability. After that the value of the second qbit is already determined!

## EPR Paradoxon

 $\frac{1}{\sqrt{2}} \left( \left| 00 \right\rangle + \left| 11 \right\rangle \right)$ 

Measuring the first qubit results in |0> or |1> with equal probability.After that the value of the second qubit is already determined!



...independent of the physical distance at the time of measurement!

⇒ No communication, no interaction can take place between the qubits (speed of light limits the distance at which communication and interaction can take place)!

This is called non-local

Physics known by then was local!

Einstein: Spooky actions at a distance



### Entanglement as Global Phenomenon



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### Entanglement: Importance

Entanglement is unique for quantum computing!

Each computation <u>not</u> involving entangled qubits, can be realized classically and **in principle** with the same efficiency than a quantum computation

(**but**: n qubits  $\Rightarrow 2^n$  classical storage | quantum parallelism | ...)

Every quantum algorithm showing exponential speedup compared to a classical algorithm, must exploit entanglement.

(R. Jozsa, N. Linden: On the role of entanglement in quantum computational speed-up. (2003) arXiv:quant-ph/0201143v2)

# Not Only Speedup

Most often, speedup is highlighted as quantum advantage

But precision can also be enhanced

• E.g. in classification





The more test data are used, the smaller is the average error in <u>classical</u> supervised learning

The average error in <u>quantum</u> supervised learning is  $Risk \ge 1 - \frac{r^2n^2 + d + 1}{d(d+1)}$ 

r: Schmidt-rank of training data n: Cardinality of S (training data) d: Dimension of Hilbert space

r = d : training data is maximal entangled  
Already n = 1 implies : Risk 
$$\ge 1 - \frac{d^2 + d + 1}{d(d+1)} \xrightarrow{d \to \infty} 0$$

A single maximal entangled element of training data suffice in high dimensions to learn with low risk a unitary transformation

# Frequency of Entanglement

Let  $\mathcal{H}$  be a Hilbert space with dim  $\mathcal{H} = d = 2^N$  and let  $\mathfrak{D}, \mathfrak{S} \subseteq \mathcal{H}$  be all states or all separable states (mixed states)

Then, 
$$\frac{\text{vol }\mathfrak{S}}{\text{vol }\mathfrak{D}}$$
 is exponentially small in N

Entanglement is ubiquitous

© Frank LeymannS. J. Szarek: Volume of separable states is super-doubly-exponentially small in the number of qubits.Phys.Rev. A72 (2005) 032304 (<a href="https://arxiv.org/abs/quant-ph/9804024v1">https://arxiv.org/abs/quant-ph/9804024v1</a>)

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#### Summary

R. Oloff: *Geometrie der Raumzeit*. Springer 2018.J. W. Robbin, D. A. Salamon: *Introduction to Differential Geometry*. Springer 2022.R. J. Adler: *General Relativity and Cosmology*. Springer 2022.

# Line Element in $\mathbb{R}^3$ (cartesian coordinates)

 $\begin{array}{c} x_{3} \\ b = (x_{1}^{b}, x_{2}^{b}, x_{3}^{b}) \\ a = (x_{1}^{a}, x_{2}^{a}, x_{3}^{a}) \\ \hline \\ x_{1} \end{array}$ Distance of two points in  $\mathbb{R}^{3}$ :  $d(a, b) = \sqrt{(x_{1}^{b} - x_{1}^{a})^{2} + (x_{2}^{b} - x_{2}^{a})^{2} + (x_{3}^{b} - x_{3}^{a})^{2}} \\ \Rightarrow d(a, b)^{2} = (x_{1}^{b} - x_{1}^{a})^{2} + (x_{2}^{b} - x_{2}^{a})^{2} + (x_{3}^{b} - x_{3}^{a})^{2} \\ = \Delta x_{1}^{2} + \Delta x_{2}^{2} + \Delta x_{3}^{2} \end{array}$ 

The so-called *line element ds* is measuring infinitesimal distances:

$$ds^2 = dx_1^2 + dx_2^2 + dx_3^2$$



The *line element ds* gives the length of infinitesimal vectors  $(dx_1, dx_2)$ :  $ds^2 = g_{11}dx_1^2 + g_{12}dx_1dx_2 + g_{22}dx_2^2$ 

The matrix  $(g_{ii})$  ist called *first fundamental form* 

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# Line Element in Curved Spaces



- M=f(U) is a curved surface (*manifold*)
- Tangent vectors in x=f(p) build the *tangent space*  $T_x M$  of M in x
- First fundamental form depends on x:  $(g_{ij}(x))$  (also:  $g_x$ )
- $g_x$  defines a scalar product on  $T_x M$ :  $\langle v, w \rangle_x \stackrel{\text{def}}{=} v^T (g_{ij}(x)) w \stackrel{\text{def}}{=} g_x(v, w)$ 
  - Reminder: a scalar product induces a metric on a vector space
  - Thus,  $g_x = (g_{ij}(x))$  is also called *Riemannian Metric* on M
  - (M, g) is called *Riemannian Manifold*



$$s(\gamma) = \int_{t_1}^{t_2} \sqrt{g_{\gamma(t)}\left(\gamma'(t), \gamma'(t)\right)} dt$$

 $s(\gamma)$  is called *arc length* of  $\gamma$  or simply *length* of the curve  $\gamma$ 

### Reminder: Directional Derivative

Let  $U \supseteq \mathbb{R}^n$ ,  $f: U \to \mathbb{R}$ ,  $F: U \to \mathbb{R}^m$  and  $v \in \mathbb{R}^n$ 

 $D_{v}f = \lim_{h \to 0} \frac{f(x + hv) - f(x)}{h} \text{ is directional derivative of f in direction of v}$ • Writing:  $D_{v}f = \partial_{v}f = \nabla_{v}f = \frac{\partial f}{\partial v}$  $D_{v}f = \operatorname{grad} f \cdot v \text{ and } D_{e_{i}}f = \frac{\partial f}{\partial x_{i}}$ 

 $D_v F (= \partial_v F = \nabla_v F)$  is build component-wise,  $D_v F = DF \cdot v \in \mathbb{R}^m$ with Jacobi-matrix  $DF \stackrel{\text{def}}{=} \left(\frac{\partial f_i}{\partial x_j}\right)_{1 \le i \le m, 1 \le i \le n}$ 

# Parallel Transport

$$X: M \to \bigcup_{p \in M} T_p M \stackrel{\text{def}}{=} TM, X(p) \in T_p M \text{ is vector field}$$

Let  $\gamma : I = [t_0, t_1] \rightarrow M$  be a curve on M. X is *parallel along*  $\gamma :\Leftrightarrow \nabla_{\gamma'(t)} X(\gamma(t)) = 0$ ( $\approx$  X doesn't change along the curve -  $\nabla$  is directional derivative)



$$\forall v_0 \in T_{\gamma(t_0)}M \exists ! V : M \to TM : V(\gamma(t_0)) = v_0 \land V \text{ parallel along } \gamma$$

 $P_{\gamma(t_0),\gamma(t_1)}: T_{\gamma(t_0)}M \to T_{\gamma(t_1)}M : v_0 \mapsto V(\gamma(t_1))$  is called *parallel transport* of  $v_0$ 

# Christoffel-Symbols

Let M be an n-dimensional manifold, i.e. dim  $T_pM = n$ ,  $\forall p \in M$ 

Let  $X_1, \ldots, X_n$  be vector fields that are a basis for  $T_pM$ ,  $\forall p \in M$ 

Then 
$$\nabla_{X_i} X_j = \sum_{k=1}^n \Gamma_{ij}^k X_k$$
, with the *Christoffel-Symbols*  $\Gamma_{ij}^k \in \mathbb{R}$ 

It is: 
$$\Gamma_{ij}^{k} = \frac{1}{2} \sum_{r} g^{kr} \left( \frac{\partial g_{jr}}{\partial x_{i}} + \frac{\partial g_{ir}}{\partial x_{j}} - \frac{\partial g_{ij}}{\partial x_{r}} \right)$$
 with  $(g^{kr}) \stackrel{\text{def}}{=} (g_{kr})^{-1}$ 

(d.h. the Christoffel-Symbols are determined by the partial derivatives of the metric)

# Directional Derivatives and Christoffel-Symbole

Let  $X_1, \ldots, X_n$  be vector fields that are a basis for  $T_p M$ 

Let  $X = \sum x_i X_i$  and  $Y = \sum y_j X_j$  be arbitrary vector fields

Then: 
$$\nabla_X Y = \sum z_k X_k$$

It is: 
$$z_k = \sum_{i,j} \Gamma_{ij}^k x_i y_j + \sum_i x_i \nabla_{X_i} y_k$$

(d.h. the directional derivatives are determined by the Christoffel-Symbols) (...and, thus, by the partial derivatives of the metric)

# Holonomy

The tangent vector v at A of curve AB is parallel transported along AB

The resulting vector at B is parallel transported along curve BC

The resulting vector at C is parallel transported along curve CA



The resulting vector w at A is in general not the original vector v!

The angle  $\alpha$  between v and w is a measure for the curvature of M

This phenomenon is referred to as *holonomy* 

 $P_{\delta(s_1),\delta(s_0)}$ 

 $P_{\gamma(t_0),\gamma(t_1)}$ 

 $\gamma(t_1), \gamma(t_0)$ 

# Riemannian Curvature

 $P_{\delta(s_0),\delta(s_1)}$ 

- Let X, Y, Z be vector fields
- $\delta, \gamma$  the curves defined by X, Y
- Z is parallel transported along
  - ... X with P<sub>\delta(s\_0),\delta(s\_1)</sub> and Y with P<sub>\gamma(t\_0),\gamma(t\_1)</sub>
    ... -X with P<sub>\delta(s\_1),\delta(s\_0)</sub> and -Y with P<sub>\gamma(t\_1),\gamma(t\_0)</sub>

$$\Rightarrow P_{\gamma(t_1),\gamma(t_0)} \circ P_{\delta(s_1),\delta(s_0)} \circ P_{\gamma(t_0),\gamma(t_1)} \circ P_{\delta(s_0),\delta(s_1)}(Z) = v$$

• Holonomy 
$$\Rightarrow v \neq Z(\delta(0)) \Rightarrow curvature!$$

• Via  $s_1 \rightarrow s_0$  and  $t_1 \rightarrow t_0$  the curves become infinitesimal and an indicator of the curvature at  $q = \delta(s_0)$  results

$$\Rightarrow \lim_{t_1 \to t_0} \lim_{s_1 \to s_0} P_{\gamma(t_1), \gamma(t_0)} \circ P_{\delta(s_1), \delta(s_0)} \circ P_{\gamma(t_0), \gamma(t_1)} \circ P_{\delta(s_0), \delta(s_1)}(Z) = \left(\nabla_X \nabla_Y - \nabla_Y \nabla_X - \nabla_{[X,Y]}\right) Z \stackrel{\text{def}}{=} R(X, Y) Z$$

No assumptions are being made about a space "embracing" the manifold M: *inner geometry*! Otherwise: *outer geometry*!

## Ricci Curvature

R(X,Y)Z defines for  $p \in M$  a multi-linear map

 $\operatorname{Rm}_p: T_pM \times T_pM \times T_pM \to T_pM$ 

as follows: for  $Y, Z \in T_p M$  fixed, define  $\Phi_p^{Y,Z} : T_p M \to T_p M$  as  $\Phi_p^{Y,Z}(X) := \operatorname{Rm}_p(X, Y, Z) = R(X, Y)Z$ 

Then,  $\operatorname{Ric}_p : T_pM \times T_pM \to \mathbb{R}, (Y, Z) \mapsto \operatorname{Tr} \Phi_p^{Y, Z}$  is called *Ricci-Map* With  $X \in T_pM$ ,  $\|X\| = 1$ ,  $\operatorname{Ric}_p(X, X)$  is the *Ricci-Curvature* in direction of X

Define 
$$R_{ij} := \sum_{a=1}^{n} \frac{\partial \Gamma_{ij}^{a}}{\partial x_{a}} - \sum_{a=1}^{n} \frac{\partial \Gamma_{ai}^{a}}{\partial x_{j}} + \sum_{a=1}^{n} \sum_{b=1}^{n} \left( \Gamma_{ab}^{a} \Gamma_{ij}^{b} - \Gamma_{ib}^{a} \Gamma_{aj}^{b} \right)$$

(d.h. the  $R_{ij}$  are defined by means of the Christoffel-Symbole and their derivatives, i.e. by the metric)

Then 
$$R := \sum_{i,j} g_{ij} R_{ij}$$
 is called *Ricci-scalar* or *scalar curvature*  
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### Geodesics



A curve  $\gamma : I \to M$  is called *geodesic* : $\Leftrightarrow \nabla_{\gamma'} \gamma' = 0$ 

• The tangent vector  $\gamma'$  does not change along the curve

• The curve has no curvature within M

Locally, a geodesic is always the shortest connection between two points

(A geodesic on a sphere is always a segment of a great-circle: the red geodesic is the shortest connection between p and q, but the green geodesic connects p and q too [but is not the shortest connection])

A curve 
$$\gamma = (\gamma_1, \dots, \gamma_n)$$
 is a geodesic  $\Leftrightarrow$  For  $1 \le i \le n$ :  

$$\frac{d^2 \gamma_i}{dt^2} + \sum_{j,k} \Gamma_{jk}^i \frac{d\gamma_j}{dt} \frac{d\gamma_k}{dt} = 0$$

(Reminder: the  $\Gamma_{ik}^{i}$  are determined by the metric and its partial derivatives!)

# Einstein's Field Equation



Matter (Stress-Energy Tensor) results in cuvature (Ricci Tensor) of space-time, such that particles move on geodetics (metric)

(a system of 16 partial differential equations of 2nd order)

# Schwarzschild-Metric

Let M be a mass, that neither rotates nor is it charged

Outside M in its nearby environment, there are no other masses

Then:  $T_{\mu\nu} = 0$  (vacuum field equation)

 $\Rightarrow$  solution (in spherical coordinates) is the so-called *Schwarzschild-Solution*:

$$ds^{2} = -c^{2} \left( 1 - \frac{2GM}{c^{2}r} \right) dt^{2} + \left( 1 - \frac{2GM}{c^{2}r} \right)^{-1} dr^{2} + r^{2} d\theta^{2} + r^{2} \sin^{2} \theta d\phi^{2}$$

Solution has two singularities at r = 0 and at  $r_S := \frac{2GM}{c^2}$  (*Schwarzschild Radius*) • In proper coordinates,  $r_S$  is no longer a singularity

• Singularity at r = 0 is a proper singularity (i.e. independent of any chosen coordinate system)

# (Anti-) de Sitter

Field equation with cosmological constant:  $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu}$ 

- $\Lambda$  has been introduced by Einstein in order to get a static universe
- <u>But</u>: in the meantime we know that the universe is not static

New observations can be explained with  $\Lambda$ 

- $\Lambda > 0 \Rightarrow$  anti-gravitation ("dark energy")  $\Rightarrow$  expansion of the universe
- $\Lambda < 0 \Rightarrow$  contraction of the universe

#### Vacuum solutions with cosmological constant:

- *de Sitter* space  $(\Lambda > 0) \Rightarrow$  constant positive curvature ( $\triangleq$  sphere)
  - …matched by observations
- Anti de Sitter space ( $\Lambda < 0$ )  $\Rightarrow$  constant negative curvature ( $\triangleq$  saddle surface)
  - Space does <u>not</u> expand
  - …does <u>not</u> match observations

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M. Camenzind: *Faszination kompakte Objekte*. Springer 2021.W. Schmitz: *Understanding Relativity*. Springer 2022.D. Grumiller, M. M. Sheikh-Jabbari: *Black Hole Physics*. Springer 2022.

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# Escape Velocity

... is the velocity  $v_F$ , which a mass m at a distance r from a mass M must have to escape from M

I.e. the kinetic energy of m must be equal to the binding energy within the gravitational field of M:

$$\frac{1}{2}mv_F^2 = \frac{GMm}{r} \quad \Rightarrow \quad v_F = \sqrt{\frac{2GM}{r}}$$

Earth:  $v_F = 11,2 \text{ km/s}$ Sun:  $v_F = 617,4 \text{ km/s}$ 

### Event Horizon

With 
$$r = r_S = \frac{2GM}{c^2}$$
 the escape velocity becomes  $v_F = \sqrt{\frac{2GM}{r_S}} = c$ !

 $\Rightarrow$  From this area even light cannot escape !

 $\Rightarrow$  This area appears to be completely black: *black hole* 

Out of the area within the Schwarzschild radius, no information at all can reach us, i.e. this area is for external observers eventless.

Thus, the sphere with radius  $r_S$  is called *event horizon* 

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# Cosmic Censorship

Penrose-Hawking Theorems prove:

Singularities are consequences of gravitational collapse, all collapsed matter will be concentrated to a single point.

Do singularities exist, which are not surrounded by an event horizon, i.e. that are directly observable (so-called *naked* singularities)?

<u>Conjecture</u> (up to now without proof):

Naked black holes do not exist.

Surface Gravity

Gravitational acceleration g at the surface r of a mass M

results from 
$$mg = \frac{GmM}{r^2}$$
 d.h.  $g = \frac{GM}{r^2}$ 

With 
$$r = r_S = \frac{2GM}{c^2}$$
 the gravitational acceleration  $\kappa_H$   
at the event horizon is:  $\kappa_H = \frac{GM}{(2GM/c^2)^2} = \frac{c^4}{4GM}$ 

 $\kappa_H$  is called *surface gravity* at the event horizon

*The surface gravity*  $\kappa_H$  *at the event horizon is constant.* 

# Horizon Theorem

Given two black holes of masses  $M_1, M_2$ , which merge

The mass of the merged black hole is then  $M = M_1 + M_2$ 

The Schwarzschild radius of the merged black hole is:

$$r_{S} = \frac{2GM}{c^{2}} = \frac{2G(M_{1} + M_{2})}{c^{2}} = \frac{2GM_{1}}{c^{2}} + \frac{2GM_{2}}{c^{2}} = r_{S,1} + r_{S,2}$$

The surface area of the horizon of the merged black hole is:

$$A_{H} = 4\pi r_{S}^{2} = 4\pi (r_{S,1} + r_{S,2})^{2} = 4\pi r_{S,1}^{2} + 4\pi r_{S,2}^{2} + 8\pi r_{S,1} r_{S,2}$$
$$= A_{H,1} + A_{H,2} + 8\pi r_{S,1} r_{S,2} > A_{H,1} + A_{H,2}$$

The horizon area  $A_H$  of black holes cannot shrink.

(experimental confirmation via measurements of gravitational waves in 2021)

# Decay of Black Holes?

<u>Assumption</u>: A black hole of mass  $M = M_1 + M_2$  decays

into two black holes of masses  $M_1, M_2$ 

As shown before:  $A_H = A_{H,1} + A_{H,2} + 8\pi r_{S,1}r_{S,2} > A_{H,1} + A_{H,2}$ 

 $\Rightarrow$  By decay, the horizon area would shrink: contradiction!

A black hole cannot decay (into black holes of smaller masses).

## No-Hair-Theorem

All properties of a black hole are completely determined by its mass M, its angular momentum J and its electric charge C.

I.e. no other physical quantities (magnetic field, number of particles, their spin,...) influence the properties of a black hole

These additional quantities ("hair") are irrelevant for black holes

Black holes have no hair

But the particles that collapsed into the black hole had these quantities! <u>Problem</u>: What happened to this corresponding information? (see later)

### Vacuum Fluctuation

Reminder: Heisenberg Uncertainty Relation  $\Delta_{v}(A) \cdot \Delta_{v}(B) \ge \frac{1}{2} \left| \left\langle [A, B] \right\rangle_{v} \right|$ 

This gives: 
$$\Delta E \Delta t \ge \frac{\hbar}{2} \text{ mit } \hbar := \frac{h}{2\pi}$$

 $E = mc^2 \Rightarrow$  for a small period of time  $\Delta t$  a mass  $\Delta m$  can come into existence

Conservation laws of physics: this matter consists of particle/antiparticle pairs

 $\approx$  Conception: in the vacuum these pairs are created permanently, but they are destroyed in a very short time ("annihilation")

"The vacuum is fluctuating"

#### HAS

### Evaporation of Black Holes

- Close to the event horizon,  $T, \overline{T}$  are created because of vacuum fluctuation
- Assumption:  $\overline{T}$  (or T) falls into the black hole, T (or  $\overline{T}$ ) escapes  $\Rightarrow$  No annihilation possible: Violation of the conservation of energy!
- <u>Analysis</u>: The energy of  $\overline{T}$  is negative ( $\Rightarrow$  conservation of energy)!

 $\Rightarrow \overline{T}$  has negative mass ( $E = mc^2$ )

 $\Rightarrow$  The black hole is loosing mass

"Evaporation of the black hole"

$$t_{\text{evaporation}} = 2.1 \times 10^{67} \left(\frac{M}{M_{\odot}}\right)^3 \text{ [years]}$$

Stellar black holes ( $M \ge M_{\odot}$ )  $\Rightarrow t_{\text{evaporation}} > 10^{67}$  [years]  $\gg 1.4 \times 10^{10}$  [years] (age of the universe)  $\Rightarrow$  Evaporation is irrelevant at present! (but: primordial black holes!)

# Macro-/Microstates

A system consists of N particles. We distinguish:

- *Microstate*: state of an individual particle of the system
- *Macrostate*: overall state of the whole system

The macrostate result from the microstates

<u>Example</u>: The macrostate of a gas (temperature, pressure,...) results from the microstates (location, momentum) of the gas particles.

*Phase space*: set of all possible microstates

Let  $p_i$  be the probability that the i-th particle is in a certain microstate

*Configuration*: particular probability distribution of the microstates (each point in the phase space has an associated probability)

# Entropy

Information that is needed to determine the configuration of a certain macrostate is called *entropy S* 

$$S := -k_B \sum_{i=1}^N p_i \cdot \ln p_i$$

The term has its origin in thermodynamics

There it is proven: 
$$dS = \frac{dQ}{T}$$
 (Q: amount of heat, T: temperature)  
Second law of thermodynamics:  $dS \ge 0$ 

# Quantum Mechanical Interpretation of Entropy

The microstate of particle *i* is the point |*i*> in Hilbert space
The phase space is substituted by the Hilbert space

The corresponding macrostate is the density matrix  $\rho = \sum p_i |i\rangle \langle i|$ 

von Neumann entropy of  $\rho$  is  $S := -\operatorname{Tr}(\rho \ln \rho)$ 

Quantum physical entropy corresponds up to the factor  $k_B$ to the information theoretical entropy

# Summary: Entropy

Entropy measures the information needed to completely describe the macrostate of a system based on the distribution of its microstates

• Entropy can never decrease

• Things with entropy have a temperature

# Bekenstein-Hawking Formula

A body with mass m and entropy S falls into a black hole

- $\Rightarrow$  Entropy outside the event horizon is reduced
  - Because: the information content of the body is no longer available outside!
- $\Rightarrow$  The mass and, thus, the horizon area of the black hole increases

Because entropy cannot decrease, the otherwise lost entropy must correspond to the increase of the horizon area

 $\Rightarrow$  The horizon area is a measure for the entropy of the black hole

The entropy of a black hole is

$$S = \frac{k_B}{4\mathfrak{L}_P^2}A_H$$

with the *Planck Length*  $\mathfrak{L}_P = \sqrt{G\hbar/c^3}$  ( $\approx 1,616 \cdot 10^{-35}$  m)

(smallest length in which space can be subdivided; everything smaller that  $\mathfrak{Q}_P$  collapses to a black hole)



# Hawking Radiation

Remember: Things with entropy have a temperature

Temperature of a black hole is  $T = \frac{\hbar c^3}{8\pi GMk_B} = \frac{\hbar}{2\pi ck_B}\kappa_H$ Note:  $T \propto \frac{1}{M}$ 

But a body with a temperature emits thermal radiation

 $\Rightarrow$  a black hole radiates: contradiction(!)

because nothing can escape a black hole!

<u>Solution</u>: this radiation corresponds to the particles, that are created by the evaporation of the black hole

 $\Rightarrow$  By radiating, the black hole loses mass, thus, it gets hotter until it explodes!

#### Holographic Principle (t'Hooft, Susskind)

Information content of a black hole proportional to the number of potential microstates within the black hole

Thus, information content should be proportional to the volume of the black hole

• ... because the particles are scattered across the volume of the black hole

But the information content is proportional to the horizon area  $S = \frac{k_B}{4\Omega_P^2} A_H$ 

- Via entanglement the number of degrees of freedom in the volume becomes proportional to its enclosing surface<sup>(\*)</sup>
- The degrees of freedom of the microstates correspond 1–1 to the degrees of freedom at the horizon : *Holographic Principle*

The horizon is like a hologram of the inner of the black hole

All information about the inner is encoded on the horizon

L. Susskind: The world as a hologram. arXiv:hep-th/9409089v2 (1994)

### Horizon Area & Information

 $S = \frac{k_B}{4} \frac{A_H}{\mathfrak{L}_B^2}$ 

 $\Rightarrow$  Entropy is the number of Planck Cells  $A_H/\mathfrak{L}_P^2$ on the horizon (up to a factor)

> Planck Cell  $\mathfrak{L}^2_P$  is the smallest area, which can carry information: Planck Cell  $\triangleq$  1 Qubit

"Information equates surface"



The same can be shown for the horizon of the whole universe

"It from Qubit"

Spacetime

Quantum Information

# Information & Gravitation $S = \frac{k_B}{4\Omega_P^2} A_H$

- The event horizon is just thus big to fit the information content of all matter that has fallen into the black hole on its horizon
- This information content determines the horizon area and, thus, the Schwarzschild radius  $r_s = \sqrt{\frac{A_H}{4\pi}} = \frac{\mathfrak{L}_P}{\sqrt{\pi k_B}}\sqrt{S}$

• This information content determines also the surface gravity  $\kappa_H = \frac{GM}{r_S^2}$ 

# Interpretation

The entropy and, thus, the information content of a black hole determines the curvature of space-time "close" to the black hole

Gravity does not exist at microscopic scales (Planck Length)

Gravity is a macroscopic effect of entropy and information

Space-time consists of smallest structures (Planck Cells), each of which carry 1 Qubit of information

The qubits of the Planck Cells are entangled

# Spacetime via Emergence

Spacetime is a phenomenon of emergence, it is no longer fundamental

Spacetime is a fabric, it emerges from entangling qubits that represent Planck cells<sup>(\*)</sup>

(but nobody knows yet what these qubits are made of)

"It from Qubit"

# Information Paradox

#### **General Relativity Theory**

Before falling into a black hole, particles have more properties than just mass, angular momentum J and its electric charge — i.e. those quantities that completely determine the properties of a black hole (No-Hair Theorem)
Like spin,...

This additional information is lost when passing the event horizon

When a black hole evaporates, this information is not recovered

#### **Quantum Physics**

The collapse is a unitary process, thus, reversible. Especially, the information lost can be recovered!

This is a contradiction!

## Additional Argument: No-Hiding Theorem

When information is lost from a subsystem  $\mathbb{H}_A$ , which is part of the composed system  $\mathbb{H}_A \otimes \mathbb{H}_B$ , this information moves to the subsystem  $\mathbb{H}_B$ .

*Information within a closed system is never lost* (≡ Law of conservation of quantum information)

Braunstein, S. L., & Pati, A. K. Quantum information cannot be completely hidden in correlations: Implications for the black hole information paradox. Phys. Rev. Lett. 98, 080502 (2007).

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#### Agenda

Entanglement

Einstein's Field Equation

Black Holes

Properties of Black Holes

ER = EPR

Summary

J. Maldacena, L. Susskind: Cool horizons for entangled black holes. arXiv:1306.0533v2 (2014) © Frank Leymann

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# Wormhole

Introduced by H. Weyl as a possible construct<sup>(\*)</sup>

- $\equiv$  Special solution of Einstein's Field Equation<sup>(\*\*)</sup>
- Space-time consists of two "sheets"
- The "sheets" are connected by a "surface of a cylinder"

#### More precise:

Let M be the space-time and  $\Omega \subseteq_{\text{compact,conn}} M$  with  $\Omega \approx [0,1] \times \Sigma$ and  $\Sigma$  be a 3-manifold with  $\partial \Sigma \approx \mathbb{S}^2$ ,

then  $\Omega$  is called a *wormhole* in M

 $\{0\} \times \partial \Sigma$  and  $\{1\} \times \partial \Sigma$  are the horizons of the two connected black holes

(\*) H. Weyl: "Feld und Materie". Annalen der Physik. 65 (14), (1921)

© Frank Leymann (\*\*) A. Einstein, N. Rosen: The Particle Problem in the General Theory of Relativity. Phys. Rev. 48, 73 (1935)



# A First Indicator

Quantum-mechanical computations allow:

- (1) State  $|\psi\rangle$  is created in a black hole A
- ② "Some" measurements are performed at the horizon of A, and the results are send to B via classical communication
- (3) This data encodes a process, that is performed at B: as a result,  $|\psi\rangle$  appears in the black hole B



 $\Rightarrow$  In order for  $|\psi\rangle$  to get from A to B, there must be a connection between A and B: a wormhole!

Note the similarity to teleportation, which requires entanglement between A and B

Wormholes and entanglement are equivalent descriptions



- QPU\_1 and QPU\_2 are entangled
- Information is submitted from QPU\_1 to QPU\_2
- The state of QPU\_1 influences what will actually be submitted
- Thus, QPU\_2 can learn about the state of QPU\_1
- Or: at the "2-end" of the wormhole one can learn about the first black hole

Computation with gravitational theory (i.e. wormholes) and computations with quantum protocols (i.e. entanglement) deliver the same results

### ER = EPR

Wormholes <sup>(*)</sup>	[ER] A. Einstein, N. Rosen: The Particle Problem in the General
	Theory of Relativity. Phys. Rev. 48, 73 (1935)

Entanglement [EPR] A. Einstein, B. Podolsky, N. Rosen: Can quantummechanical description of physical reality be considered complete? Phys. Rev. 47, 777 (1935)

(\*) Wormholes are also called *Einstein-Rosen-Bridges* 

#### ER = EPR:

Wormholes and entanglement are the same phenomenon

J. Maldacena, L. Susskind: Cool horizons for entangled black holes. arXiv:1306.0533v2 (2013) J. Maldacena et al.: Diving into traversable wormholes. Fortschr. Phys., 65(5), 1700034 (2017)

#### <u>Consequence</u>: Entanglement between particles origin from wormholes!

J.C. Baez, J. Vicary: Wormholes and Entanglement. Classical and Quantum Gravity, Vol. 31 No. 21 (2014) arXiv:1401.3416v2

# Generation of Wormholes in Principle

- Create a huge set of Bell-pairs
- Split the entangle particles into two sets and move them apart in large distance
- Then, collapse each of these particle clouds into a black hole

 $\Rightarrow$  Two entangled black holes!



# Hawking Radiation & Wormholes

- Black holes resulting from a "simple" collapse are unilateral, i.e. they are not *mouth* of a wormhole
- Particles resulting from evaporation of a black hole are entangled with the inner black hole
  - Image: ...and, thus, connected with the black hole by means of "small" wormholes
- By collapse of (a subset of) the emitted particles a "big" wormhole with two mouths results
  - …time required for this process is called Page-Time

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### Wormhole Traversability<sup>(\*)</sup>

#### A wormhole cannot be traversed

- A wormhole is growing faster than matter can move through it<sup>(\*)</sup>
- This can be explained via complexity theory (complexity of a circuit needed to create the quantum state of the black hole)



But one can "meet in the middle"

Microscopic wormholes could be traversable

…and submit information

### Experiment: Traversable Wormhole<sup>(\*)</sup>

Based on *holographic duality*, i.e.:

- Dynamics of quantum systems ~ Effects of quantum gravity
   With QPUs one can run experiments about quantum gravity<sup>(\*\*)</sup>
- This enables quantum circuits whose execution correspond to the dynamics of traversable wormholes<sup>(\*)</sup>
  - Effects of negative energy  $\equiv$  Operations on entangled qubits

This is consistent with ER = EPR !

• This can be understood as a kind of quantum teleportation

(\*) D. Jafferis et al.: Traversable wormhole dynamics on a quantum processor. Nature volume 612, pages 51–55 (2022)
 [BLOG Summary 1: <u>https://ai.googleblog.com/2022/11/making-traversable-wormhole-with.html]</u>
 [BLOG Summary 2: <u>https://inqnet.caltech.edu/wormhole2022/]</u>

(\*\*) S. Nezami et al.: Quantum Gravity in the Lab: Teleportation by Size and Traversable Wormholes (2021)
 Part 1 arXiv:1911.06314v2 — Part II arXiv:2102.01064v1



- ③ Creation of two entangled qubits (ERP pair) at A (probe and reference)
- (4) SWAP the probe with a qubit in B ( $\triangleq$  probe in wormhole)
- (5) Circuit with chaotic evolution of the wormhole (probe gets "scrambled")
- (6) Entanglement operation between A and B ( $\triangleq$  negative energy  $\Rightarrow$  traversability)
- (7) Another circuit with chaotic evolution of the wormhole (moves probe to B)
- B Determine degree of entanglement between probe<sub>B</sub> in B and reference<sub>A</sub> in A
  - Degree of entanglement has increased  $\Rightarrow$  information has been submitted!

Within holographic duality this means: The probe is a particle, which has been moved from A to B

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Entanglement

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Summary

# Summary

- Einstein's Field Equation describes the inner geometry of space-time
- Black holes result as a solution of this equation
  - ...and are confirmed by observation
- Black holes have fundamental properties
  - Horizon Theorem, No-Hair Theorem, Haking-Radiation,...
- Entropy is proportional to the area of the event horizon
   ⇒ Holographic Principle
- Wormholes are a special solution of the field equation
  - ...and can be described as entanglement between black holes
- NISQ enables experiments about quantum gravity
- Deep understanding of...
  - Entanglement  $\Rightarrow$  ER = EPR
  - Spacetime (emergence from Planck-Cells)  $\Rightarrow$  *It from Qubit*

AAS

#### End